An added mass partitioned algorithm for rigid bodies and incompressible flows

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Collaborators

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<u>Support</u>

Department of Energy Office of Advanced Scientific Computing Research Applied Mathematics Program NNSA Rensselaer Polytechnic Institute In recent work we have focused on large deformation and/or displacement FSI using partitioned solvers



- Component solvers remain independent
 - can use existing solvers
 - no need to solve (or precondition) a coupled implicit system
- Can naturally take advantage of disparate time scales
 - e.g. mixing implicit and explicit integration
- High levels of algorithmic concurrency
 - maps well to modern and emerging computers

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<u>Traditional partitioned schemes have suffered from "added-mass instabilities"</u> <u>for solids which are sufficiently light when compared to the fluid</u>

- Traditional partitioned FSI algorithms (Cirak, et. al. 2007, Bungartz and Schafer 2006)
 - 1. advance fluid (using interface velocity/displacement from the solid)
 - 2. advance solid (apply fluid forces to the solid)
 - 3. possibly iterate with under-relaxation to convergence



- Some analysis of added-mass instabilities can be found in the literature, for example
 - Causin, Grebeau, and Nobile, 2005 (stability with relaxation)
 - Gretarsson, Kwatra, and Fedkiw 2011 (semi-monolithic formulations)





















Body must displace and entrain fluid to move and therefore appears more massive than in vacuum ... the so called "added mass"

As a concrete motivating example consider a rising rigid body in counterflow

- Incompressible Navier-Stokes
- Light rigid body



This case has both strong added-mass and added-damping effects



- Added mass relates to increased apparent mass owing to fluid displacement
 - i.e. the local geometry occupied by solid changes
- Added damping relates to increased apparent inertia owing to viscous fluid drag
 - i.e. the local geometry remains fixed

To understand these effects in isolation we derive extremely simple models by localizing and linearizing the problem near the interface



- y-translations relate to added-mass effects
- x-translations relate to added-damping effects

The resulting model is used to motivate our new AMP algorithms and discuss the performance of traditional partitioned (TP) schemes



$$\begin{array}{ll} \text{Fluid:} & \rho \frac{\partial \mathbf{v}}{\partial t} + \nabla p = \mu \Delta \mathbf{v}, & \mathbf{x} \in \Omega, \\ & \nabla \cdot \mathbf{v} = 0, & \mathbf{x} \in \Omega, \end{array} \\ \text{Rigid body:} & m_b \, a_u = \int_0^L \mu \frac{\partial u}{\partial y}(x,0,t) \, dx + g_u(t), \\ & m_b \, a_v = -\int_0^L p(x,0,t) \, dx + g_v(t), \end{array} \\ \text{Interface:} & \mathbf{v}(x,0,t) = \mathbf{v}_b(t), & x \in [0,L], \end{array}$$

An added-mass model problem is derived by considering vertical motions



$$\begin{cases} \rho \frac{\partial v}{\partial t} + \frac{\partial p}{\partial y} = 0, & y \in (0, H), \\ \frac{\partial v}{\partial y} = 0, & y \in (0, H), \\ m_b \frac{d v_b}{d t} = -\int_0^L p \, dx, \\ v(0, t) = v_b, \quad p(H, t) = p_H(t), \end{cases}$$

• The AMP scheme matches the vertical accelerations at the interface

$$\left. \frac{\partial v}{\partial t} \right|_{y=0} = \frac{dv_b}{dt} = a$$

And applies a generalized Robin condition to the fluid pressure equation

$$\rho a + \frac{\partial p}{\partial y} = 0 \qquad \qquad m_b a = \int_0^L p \, dx$$

The resulting AMP scheme is stable for any finite mass, while the traditional scheme suffers



$$\begin{cases} \rho \frac{\partial v}{\partial t} + \frac{\partial p}{\partial y} = 0, & y \in (0, H), \\ \frac{\partial v}{\partial y} = 0, & y \in (0, H), \\ m_b \frac{dv_b}{dt} = -\int_0^L p \, dx, \\ v(0, t) = v_b, \quad p(H, t) = p_H(t), \end{cases}$$

• The added mass for this problem is easily identified as

$$M_a = \rho L H$$

• Thm: The 2nd order accurate AMP scheme for the added-mass model problem is stable provided $m_b + M_a$ is bounded away from zero.

• Thm: The 2nd order traditional partitioned scheme is stable if and only if $m_b > M_a$

An added-damping model problem is derived by considering horizontal motions



• The AMP scheme uses the exact solution to form the discrete approximation

$$\mu \int_0^L \frac{\partial u}{\partial y}(0,t) \, dx \approx -\beta \mathcal{D} u_b$$

$$\mathcal{D} \approx \mu \int_0^L \frac{1 - e^{-\delta}}{\Delta y} \qquad \qquad \delta = \frac{\Delta y}{\sqrt{\nu \Delta t/2}}$$

here δ is a ratio of viscous length scales, and D is an added-damping coefficient (in 3D these become added-damping tensors)

<u>The AMP scheme with extra velocity projection is stable even for</u> <u>massless bodies</u>



The AMP-RB scheme is implemented in Overture and is found to be stable against both added-mass and added damping instabilities without iteration

$$\rho_b = .001$$







A more challenging case is that of a light cylinder rising in counterflow

- The AMP-RB scheme is again stable without any iteration
- Traditional partitioned scheme needs ~2000 sub iterations to stabilize



$$\rho_b = .001$$

A more challenging case is that of a light cylinder rising in counterflow

$$\rho_b = .001$$

Extensions to multiple bodies presents no particular challenges, and the scheme is robust even in the presence of both light and heavy bodies



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<u>Summary</u>

• Careful analysis of simplified model problems motivates our stable and 2nd order accurate AMP-RB scheme for incompressible flows

• Stability analysis in simple geometries shows excellent stability properties even in the uniterated form

• Implementation within Overture illustrates the utility of the approach for both light and heavy bodies

Future Work

- implement in 3D
- investigate the alternate added-mass potential formulation
- New FSI regimes (incompressible/incompressible, compressible/beams)